



# Probing Gauge Boson Quartic Couplings in Multi-boson Events

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(contributions from J Damgov, J Faulkner, Q Li, P Teles, D Yang)



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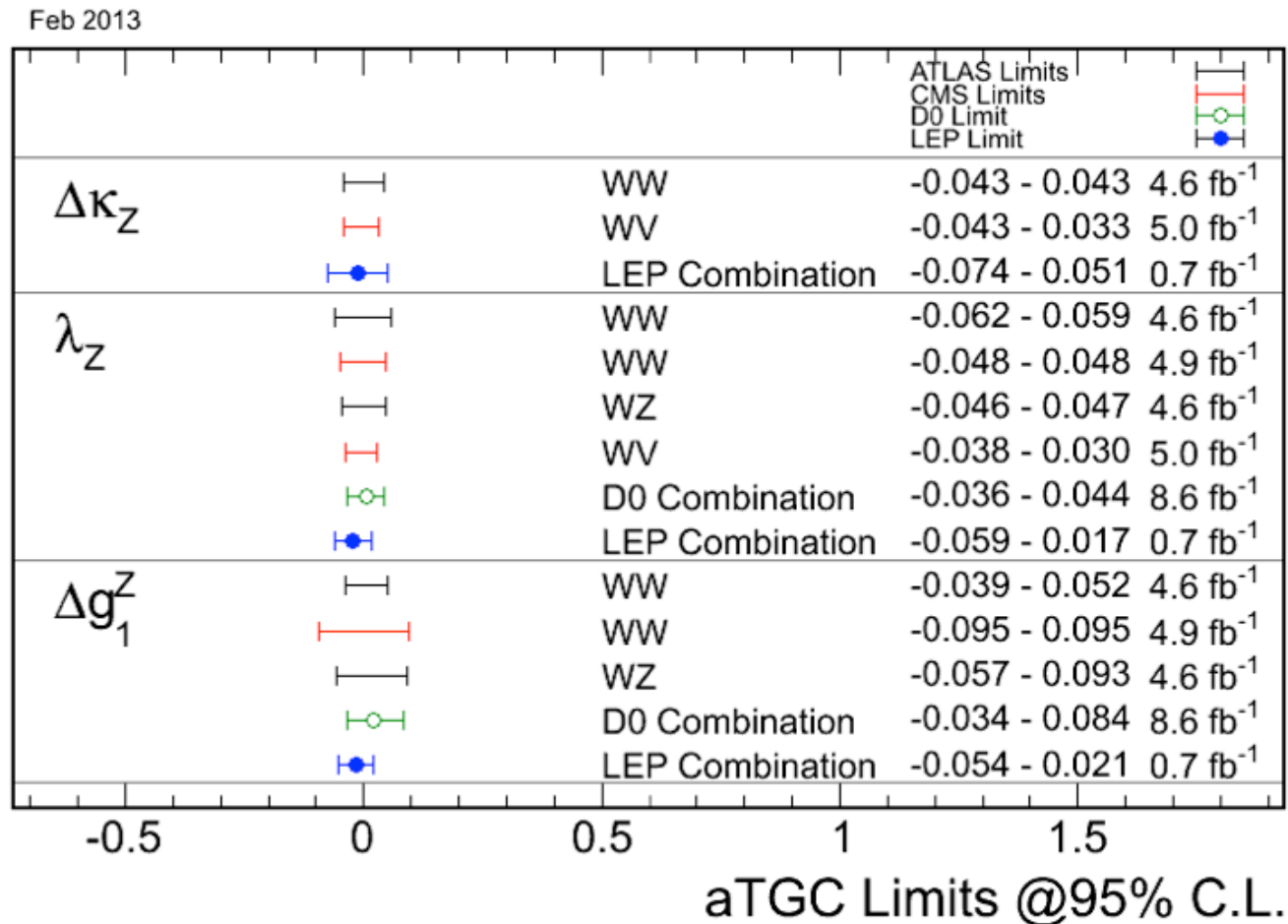
# Quartic couplings

I will only talk about couplings involving W boson

- In the SM, the allowed couplings are:  
 $WW\gamma\gamma$ ,  $WWZ\gamma$ ,  $WWWW$ ,  $WWZZ$
- Observable in two topologies at the LHC
  - Triple gauge boson production (e.g.,  $W\gamma\gamma$ ,  $WW\gamma$ ,  $WWW$ ,  $WWZ$ )
  - Scattering process (e.g.,  $\gamma\gamma \rightarrow WW$ ,  $WW \rightarrow WW$ )
- Anomalous couplings introduced via effective Lagrangian
  - Should use the linear realization with light Higgs
  - aQGCs for SM allowed processes introduced at dimension 6
  - However they are the same operators as the aTGCs which are better measured (see next slide)
- Lowest independent aQGC interactions are dimension 8

# Summary of charged aTGC measurements

In the notation of LEP parametrization [hep-ph/9601233](http://hep-ph/9601233)



- aTGCs entangled with aQGC, as explained in the following slides.

- Current constraints on aTGCs: **< 10% deviation from SM.**

- Expect to achieve a **few % precision** with 8 TeV data.

# Anomalous quartic couplings in dimension 8

All D8 aQGC operators  
in Eboli's notation

hep-ph/0606118  
Eboli et. al.

$$\mathcal{L}_{S,0} = [(D_\mu \Phi)^\dagger D_\nu \Phi] \times [(D^\mu \Phi)^\dagger D^\nu \Phi]$$

$$\mathcal{L}_{S,1} = [(D_\mu \Phi)^\dagger D^\mu \Phi] \times [(D_\nu \Phi)^\dagger D^\nu \Phi]$$

$$\mathcal{L}_{M,0} = \text{Tr} [\hat{W}_{\mu\nu} \hat{W}^{\mu\nu}] \times [(D_\beta \Phi)^\dagger D^\beta \Phi]$$

$$\mathcal{L}_{M,1} = \text{Tr} [\hat{W}_{\mu\nu} \hat{W}^{\nu\beta}] \times [(D_\beta \Phi)^\dagger D^\mu \Phi]$$

$$\mathcal{L}_{M,2} = [B_{\mu\nu} B^{\mu\nu}] \times [(D_\beta \Phi)^\dagger D^\beta \Phi]$$

$$\mathcal{L}_{M,3} = [B_{\mu\nu} B^{\nu\beta}] \times [(D_\beta \Phi)^\dagger D^\mu \Phi]$$

$$\mathcal{L}_{M,4} = [(D_\mu \Phi)^\dagger \hat{W}_{\beta\nu} D^\mu \Phi] \times B^{\beta\nu}$$

$$\mathcal{L}_{M,5} = [(D_\mu \Phi)^\dagger \hat{W}_{\beta\nu} D^\nu \Phi] \times B^{\beta\mu}$$

$$\mathcal{L}_{M,6} = [(D_\mu \Phi)^\dagger \hat{W}_{\beta\nu} \hat{W}^{\beta\nu} D^\mu \Phi]$$

$$\mathcal{L}_{M,7} = [(D_\mu \Phi)^\dagger \hat{W}_{\beta\nu} \hat{W}^{\beta\mu} D^\nu \Phi]$$

$$\mathcal{L}_{T,0} = \text{Tr} [\hat{W}_{\mu\nu} \hat{W}^{\mu\nu}] \times \text{Tr} [\hat{W}_{\alpha\beta} \hat{W}^{\alpha\beta}]$$

$$\mathcal{L}_{T,1} = \text{Tr} [\hat{W}_{\alpha\nu} \hat{W}^{\mu\beta}] \times \text{Tr} [\hat{W}_{\mu\beta} \hat{W}^{\alpha\nu}]$$

$$\mathcal{L}_{T,2} = \text{Tr} [\hat{W}_{\alpha\mu} \hat{W}^{\mu\beta}] \times \text{Tr} [\hat{W}_{\beta\nu} \hat{W}^{\nu\alpha}]$$

$$\mathcal{L}_{T,5} = \text{Tr} [\hat{W}_{\mu\nu} \hat{W}^{\mu\nu}] \times B_{\alpha\beta} B^{\alpha\beta}$$

$$\mathcal{L}_{T,6} = \text{Tr} [\hat{W}_{\alpha\nu} \hat{W}^{\mu\beta}] \times B_{\mu\beta} B^{\alpha\nu}$$

$$\mathcal{L}_{T,7} = \text{Tr} [\hat{W}_{\alpha\mu} \hat{W}^{\mu\beta}] \times B_{\beta\nu} B^{\nu\alpha}$$

$$\mathcal{L}_{T,8} = B_{\mu\nu} B^{\mu\nu} B_{\alpha\beta} B^{\alpha\beta}$$

$$\mathcal{L}_{T,9} = B_{\alpha\mu} B^{\mu\beta} B_{\beta\nu} B^{\nu\alpha}$$

$\mathcal{L}_M$  have D6  
equivalents  
( $a_0, a_c$ ),  
 $\mathcal{L}_T$  are  
novel to D8

|  | WWWW | WWZZ | ZZZZ | WWAZ | WWAA | ZZZA | ZZAA | ZAAA | AAAA |
|--|------|------|------|------|------|------|------|------|------|
| $\mathcal{L}_{S,0}, \mathcal{L}_{S,1}$                                       | X    | X    | X    | O    | O    | O    | O    | O    | O    |
| $\mathcal{L}_{M,0}, \mathcal{L}_{M,1}, \mathcal{L}_{M,6}, \mathcal{L}_{M,7}$ | X    | X    | X    | X    | X    | X    | X    | O    | O    |
| $\mathcal{L}_{M,2}, \mathcal{L}_{M,3}, \mathcal{L}_{M,4}, \mathcal{L}_{M,5}$ | O    | X    | X    | X    | X    | X    | X    | O    | O    |
| $\mathcal{L}_{T,0}, \mathcal{L}_{T,1}, \mathcal{L}_{T,2}$                    | X    | X    | X    | X    | X    | X    | X    | X    | X    |
| $\mathcal{L}_{T,5}, \mathcal{L}_{T,6}, \mathcal{L}_{T,7}$                    | O    | X    | X    | X    | X    | X    | X    | X    | X    |
| $\mathcal{L}_{T,8}, \mathcal{L}_{T,9}$                                       | O    | O    | X    | O    | O    | X    | X    | X    | X    |

# aQGC D6 vs D8

- In the two realizations
  - Linear: all lowest order independent aQGCs are D8
  - Nonlinear: a number of dimensions, aQGCs involving  $\gamma$  are D6
- Consider  $WW_{\gamma\gamma}$ 
  - Largest contributing nonlinear terms:
    - Limits set on  $a/\Lambda^2$ 

$$L_6^0 = -\frac{e^2}{16\Lambda^2} a_0 F^{\mu\nu} F_{\mu\nu} \vec{W}^\alpha \cdot \vec{W}_\alpha$$

$$L_6^c = -\frac{e^2}{16\Lambda^2} a_c F^{\mu\alpha} F_{\mu\beta} \vec{W}^\beta \cdot \vec{W}_\alpha$$
    - Equivalent D8 terms ( $L_{M2}, L_{M3}$ )
      - Limits set on  $q/\Lambda^4$
      - Straightforward conversions
 
$$\frac{q_i}{\Lambda^4} = \frac{8a_i}{\Lambda^2 M_W^2}$$
- Expectations:
  - SM rate detectable with TGC and QGC contributions at  $e^2$
  - aTGC and aQGC entangled, suppressed by  $q/\Lambda^4$
  - Sensitivity on high  $p_T$  tail

## Burden of legacy ....

**Almost all previous work in nonlinear realization**

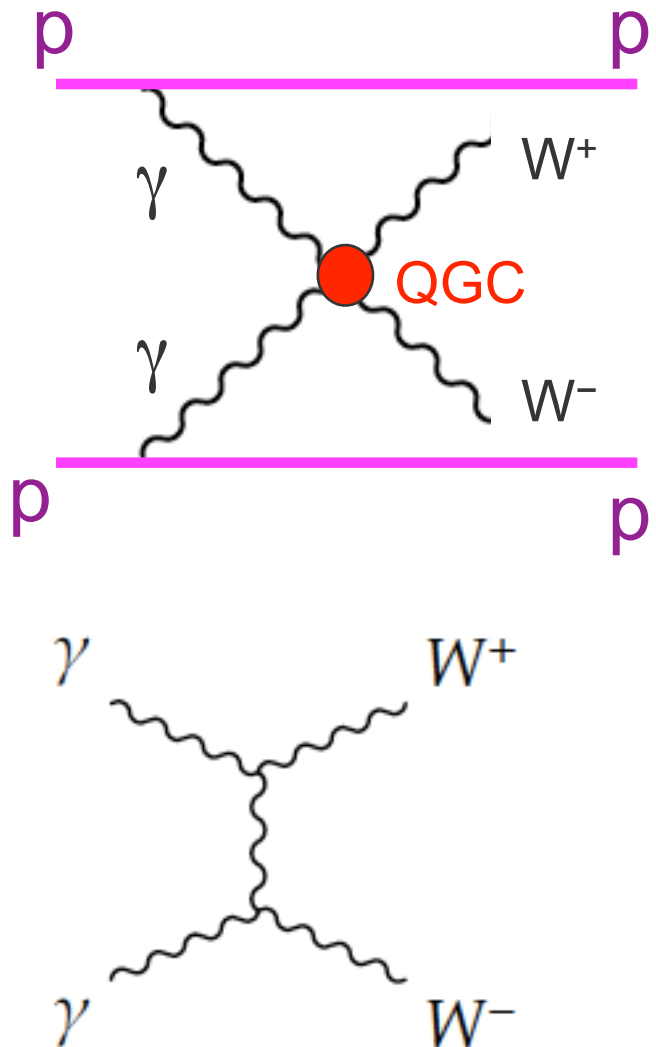


- Symmetries enforced without light Higgs
- Lower dimension D4, D6 aQGCs
- Have to connect with that work
  - LEP, LHC limits already set in that approach
  - they often use arbitrary form factors to dampen non-unitarity

### **Proposal to manage this burden**

- Adopt D8 (linear) approach for setting aQGC limits
- However, in order to easily compare with the existing results
  - use D6 equivalents for operators that exist in both D6 and D8
- Operators that are novel in D8 are probed for the first time, so there is no legacy issues to take care of

# Quartic couplings in $\gamma\gamma \rightarrow WW$ process



arXiv:1305.5596

See talk by Jonathan Hollar in LHC-EWK-WG meeting on April 16.

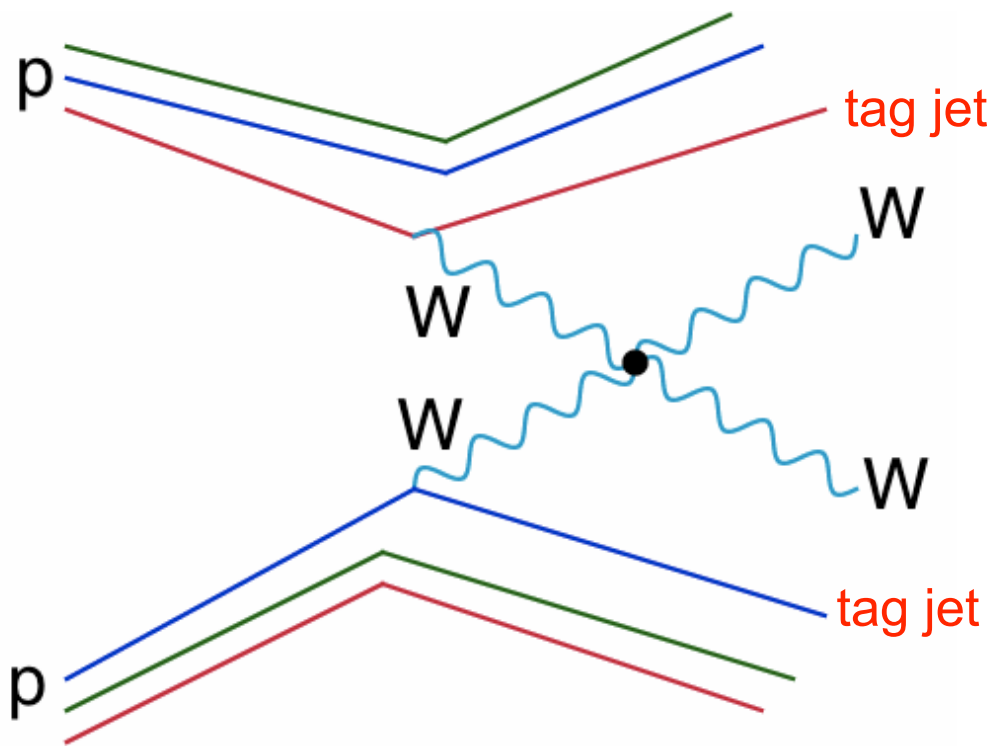
Limits on aQGC without form-factors:

$$-2.80 \times 10^{-6} < a_0^W / \Lambda^2 < 2.80 \times 10^{-6} \text{ GeV}^{-2}$$

$$-1.02 \times 10^{-5} < a_C^W / \Lambda^2 < 1.02 \times 10^{-5} \text{ GeV}^{-2}$$

$\mathcal{O}(10^2)$  times more constraining than the LEP combined limit

# Quartic couplings via $WW \rightarrow WW$ VBF



arXiv:1303.6335, 1304.4599,  
1304.0080, 1305.0251,  
1212.4158, 1212.3598,  
1209.2389, 1203.2771,  
1112.1171, hep-ph/  
0201098, ....

**Lot of theoretical interest,  
only way to measure  
WWWW coupling,  
important for  
understanding EWSB.**

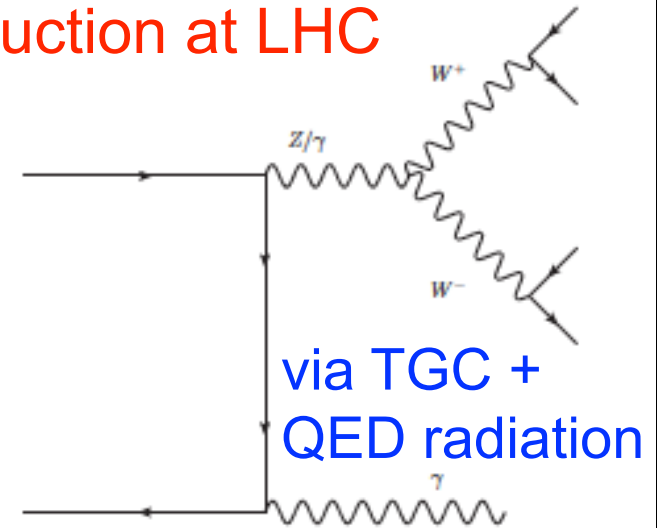
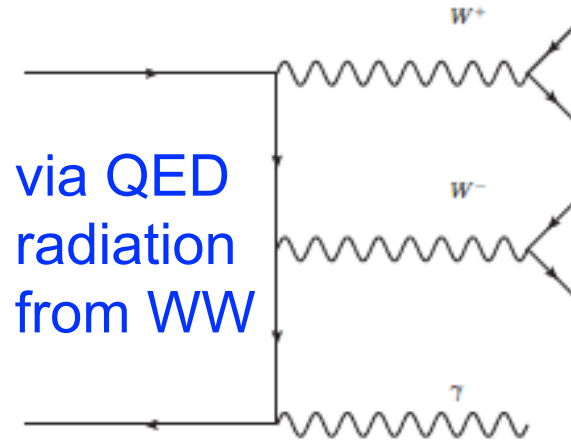
**Experimentally much harder, no measurement yet.**

# Probing quartic couplings via VV production

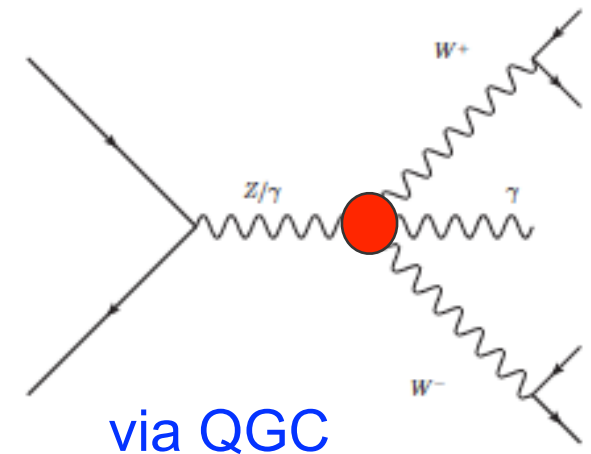
For example:  $WW\gamma$  production at LHC

References:

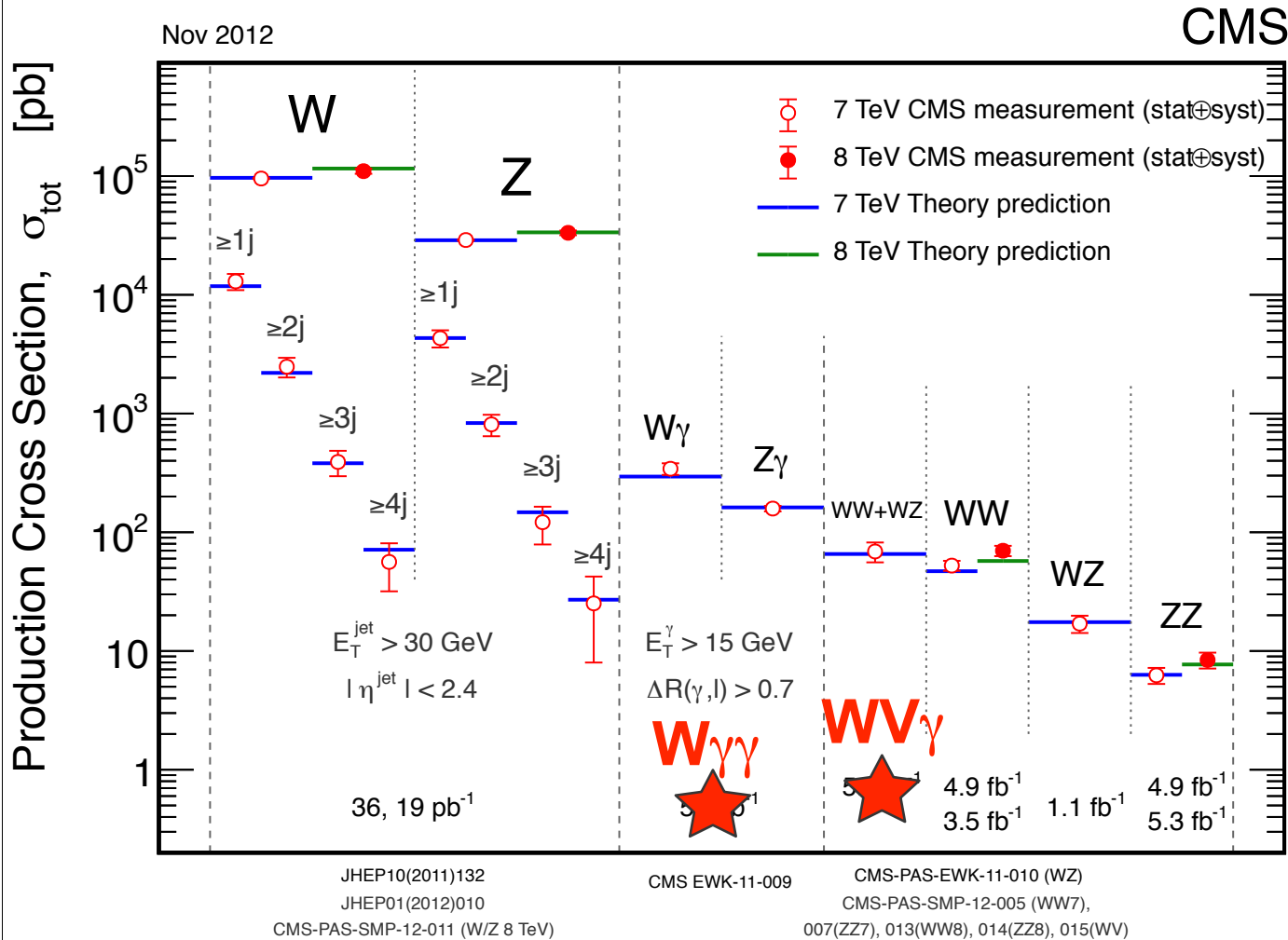
- 1.) Yang et al, arXiv: 1211.1641
- 2.) LEP combination, hep-ex/0612034
- 3.) Bozzi et al, arXiv: 0911.0438



- SM production highly suppressed
  - By a factor of  $10^3$  compared to WW
- aQGC at  $WW\gamma\gamma$  and  $WW\gamma Z$  vertices can enhance production for high photon  $p_T$  events by several factors



# Truly rare processes: sub-pb cross section



- Higher BR makes semi-leptonic channel attractive

$-\sigma \times \text{BR for } W\bar{V}\gamma \approx 60 \text{ fb}$

w/o cut on photon  $p_T$ , where  $V = W$  or  $Z \rightarrow qq$

## WW $\gamma$ , WZ $\gamma$ semi-leptonic channel expectations

- Within detector fiducial, expect 10–20 reconstructed WW $\gamma$  events ( $\gamma + \ell + E_T^{\text{miss}} + jj$ ) in 20 fb $^{-1}$  of 8 TeV data
- Given small S/B, barely getting sensitive to SM WW $\gamma$  signal  
– likely to set upper limit @ a few times the SM cross section
- Expect more constraining limits on aQGC than LEP

### Simulation

LO Madgraph simulation

- process: p p  $\rightarrow$  w+ w- a @ 8TeV LHC
- PDF (LO): CTEQ6L1, scale: default MadGraph setting
- generator cuts:  $p_T^\gamma > 10$  GeV,  $|\eta_\gamma| < 2.5$ ,  $\Delta R(\gamma, j) > 0.5$   
(not Rja cut, but the cut as Eq.(3.4) in [arXiv:0911.0438](https://arxiv.org/abs/0911.0438))

$$\sum_{i, R_{i\gamma} < R} p_T^{\text{parton}, i} \leq \frac{1 - \cos R}{1 - \cos \delta_0} p_T^\gamma \quad \forall R \leq \delta_0,$$

# NLO simulation & computation of k-factors

<http://amcatnlo.cern.ch/>

NLO QCD matched with Parton Shower (HERWIG or PYTHIA)

generate p p > w+ w- a [QCD]

output nlowwa

launch -m

4 core mode on a single 3.3GHz machine,

**~21 hours to get 40k events**

Output: (1) **events.lhe.gz** unweighted events (up to a sign),

NLO matched with Parton shower level

(2) **events\_HERWIG6\_0.hep.gz** stdHEP file, showered events

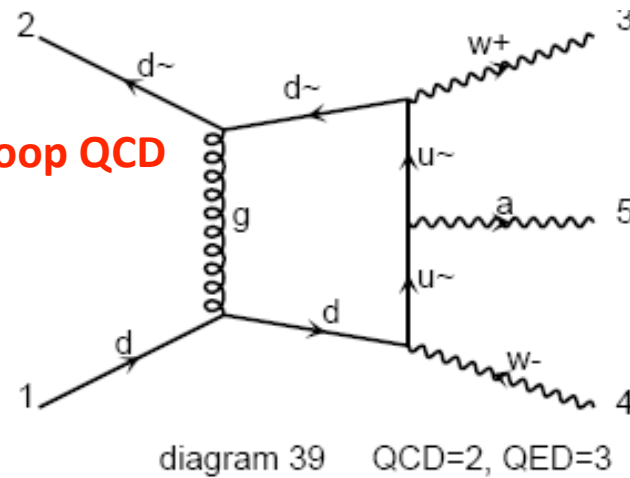
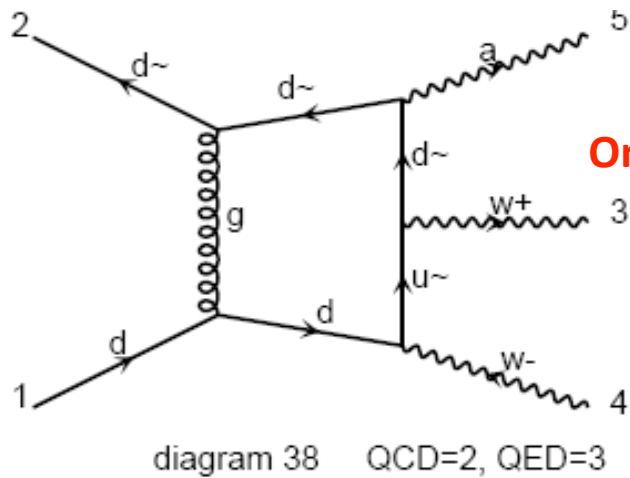
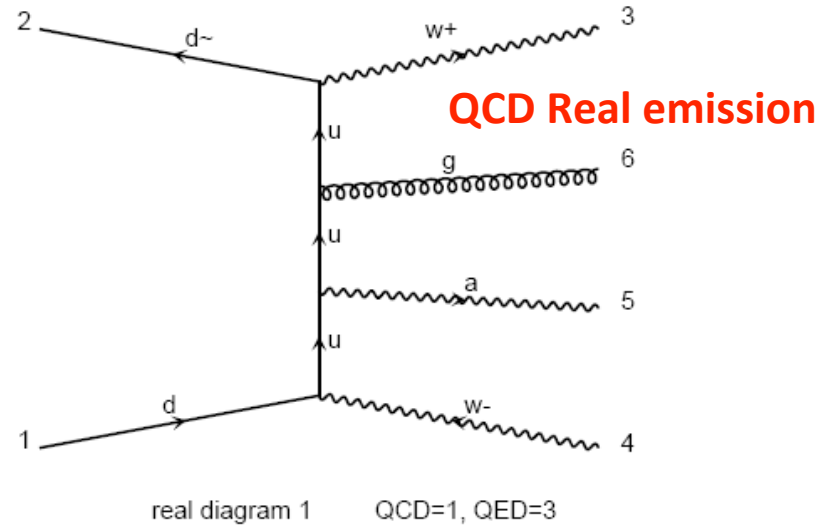
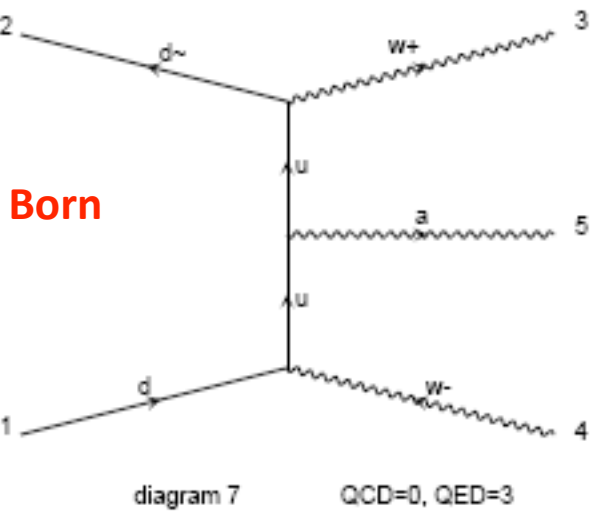
Total cross sections:

LO:  $0.1428 \pm 0.0002$  pb

NLO (CTEQ6M PDF):  $0.2533 \pm 0.0011$  pb

**K factor: 1.8**

# Some representative diagrams from aMC@NLO



## Event selection and expected yields

### ▶ Event selection:

- Lepton  $p_T > 25$  GeV,  $|\eta| < 2.4$ , MET  $> 35$  GeV
- At least 2 non-b jets with  $p_T > 30$  GeV,  $|\eta| < 2.5$
- Photon  $E_T > 30$  GeV,  $|\eta| < 1.44$ ,  $\Delta R(\gamma, \ell) > 0.5$ ,  $\Delta R(\gamma, j) > 0.5$
- $|\Delta\eta(j1, j2)| < 1.4$
- $70 < M_{jj} < 100$  GeV for the leading central jets

▶ Expected yields in 8 TeV,  $20 \text{ fb}^{-1}$  data with optimized selection:  
340 events, 12  $WV\gamma$  signal and 328 background ( $W\gamma$ +jets,  $WV$ +fake photon,  $t\bar{t} + \gamma$ , multi-jet)

➔ **NOT significant enough to see SM production**

▶ Use  $\gamma p_T$  as observable for setting limits on aQGC.

# $WV\gamma$ k-factor depends on photon $p_T$ ☹️

Also different for SM and aQGC. Specially crucial to account for this difference at high  $p_T$ .

At low  $p_T$  (<100 GeV):

- SM & aQGC have ~same k-factor

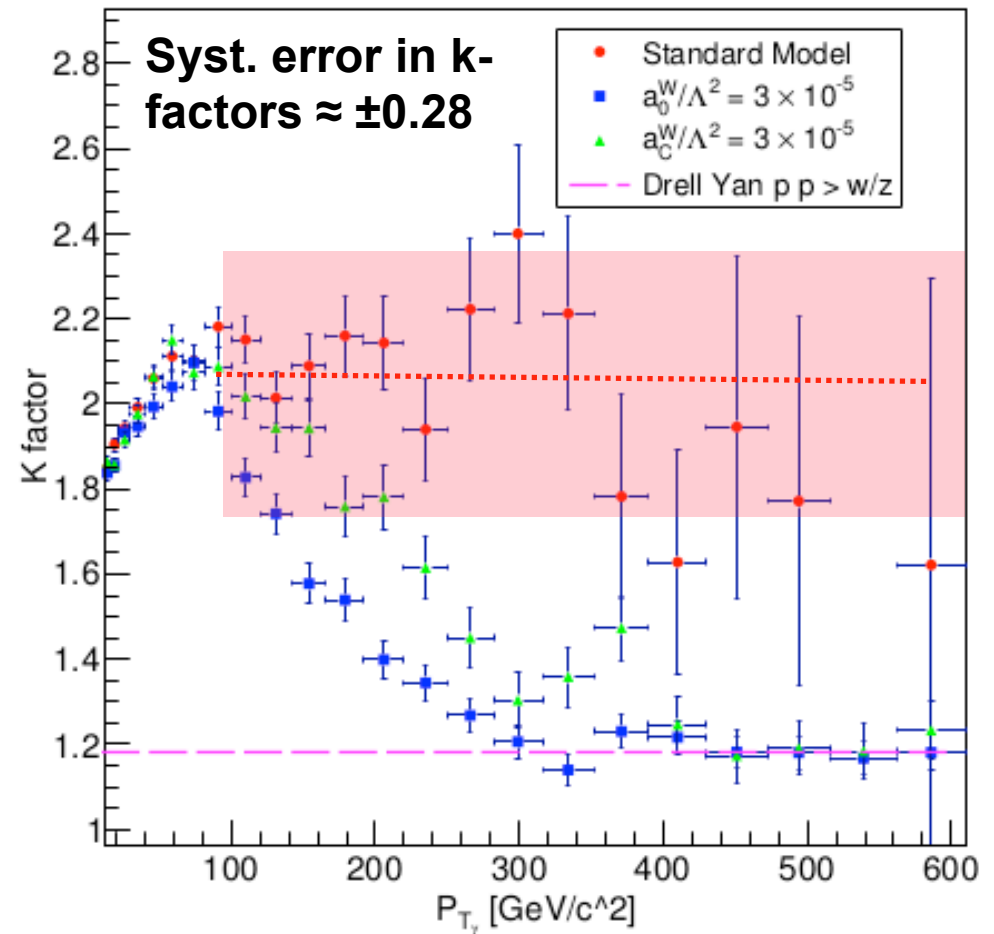
At high  $p_T$  (>300 GeV):

- SM (Red) k-factor ~2
- aQGC behaves like Drell-Yan (magenta): k-factor ~1.185

The near-constant k-factor at large  $p_T$  can be explained by the fact that only one Feynman diagram (4-gauge boson vertex one) get enhanced by the aQGC.

At medium  $p_T$  (100–300 GeV):

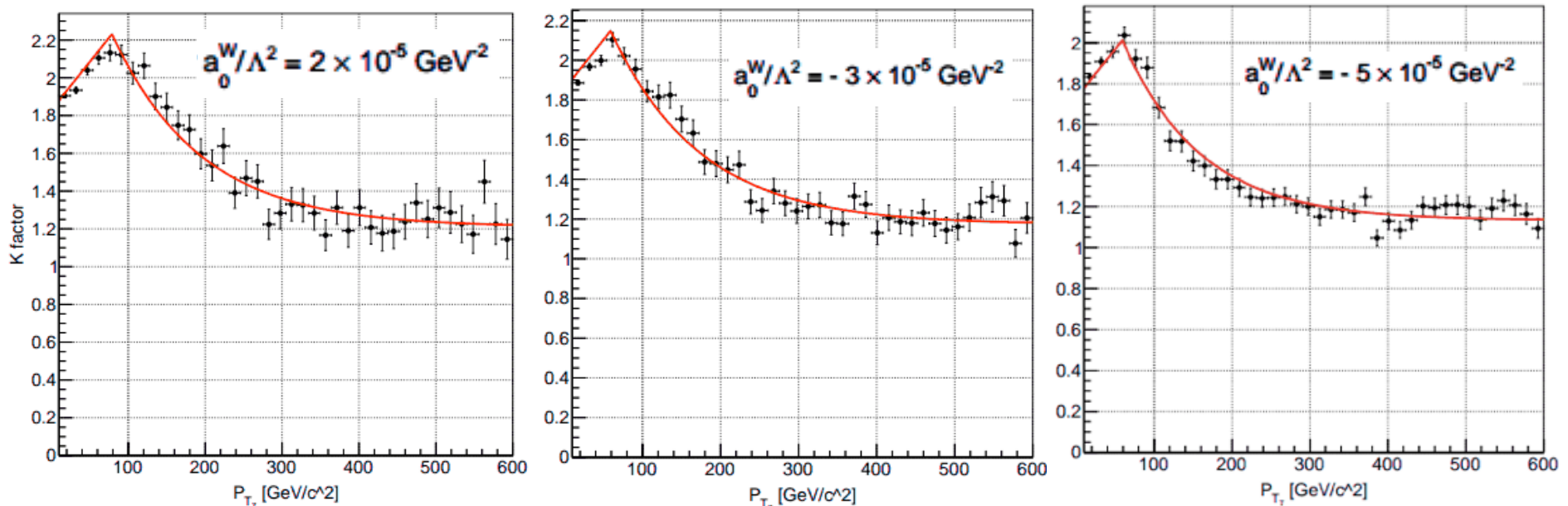
- aQGC k-factor decreases exponentially



# Attempt to parametrize aQGC k-factors

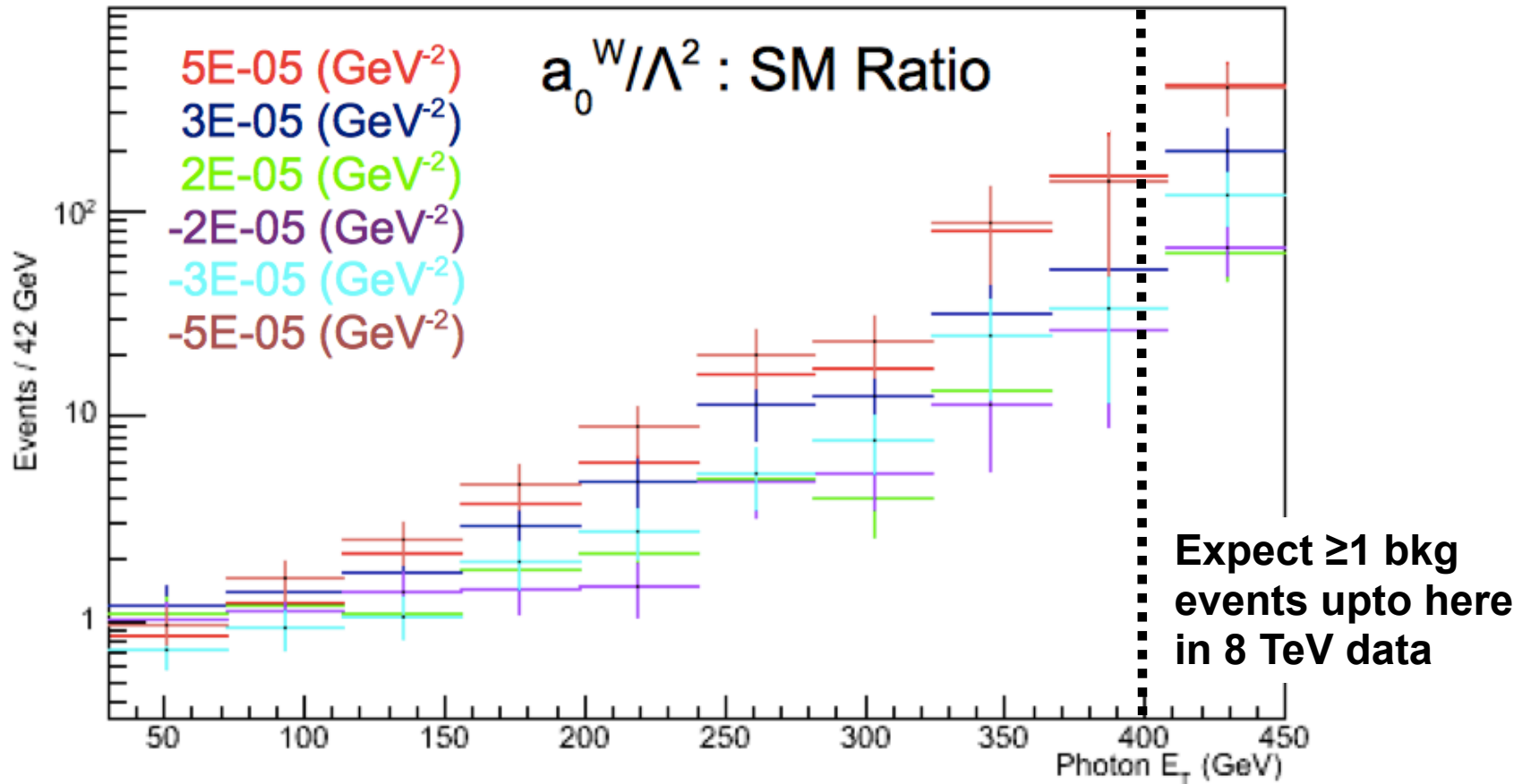
**Parametrization:** Linear below the turnover point, exponential above it.

$$\begin{aligned} & (a - 1.185) * \exp(-b * (x - c)) + 1.185 && \text{for } x > p_{T0} \\ & d * x + (a - d * c) && \text{for } x < p_{T0} \end{aligned}$$



Although parametrization can be improved, limits derived using discrete k-factor and smooth function are essentially identical.

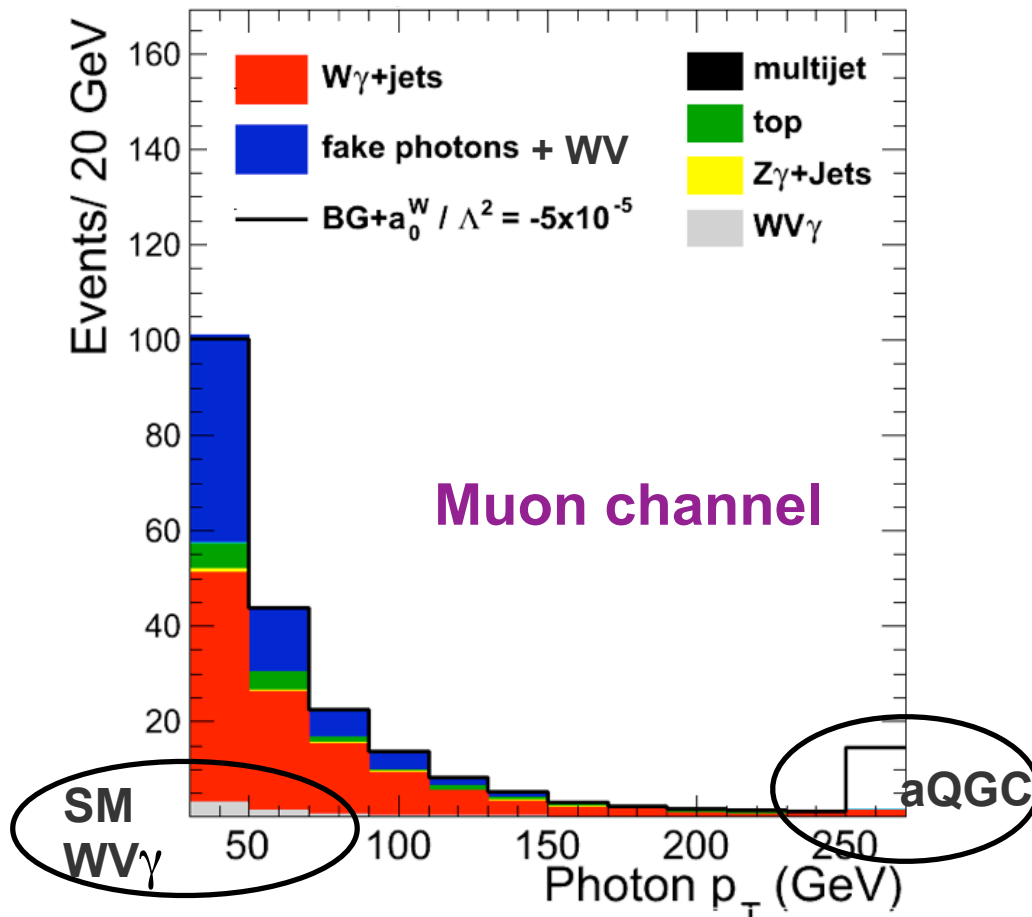
# Sensitivity to aQGC



Studies ongoing to determine unitarity band, i.e., values of  $\gamma$   $p_T$  and aQGC combination for which prediction becomes non-unitary.

# Observable: $\gamma$ $p_T$ distribution

$$\int L dt = 20 \text{ fb}^{-1}$$



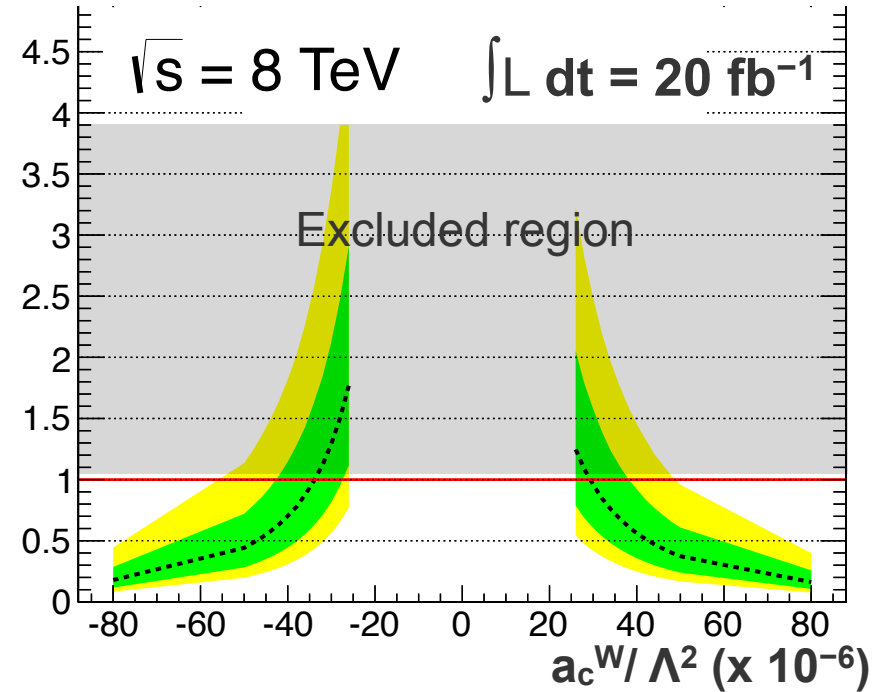
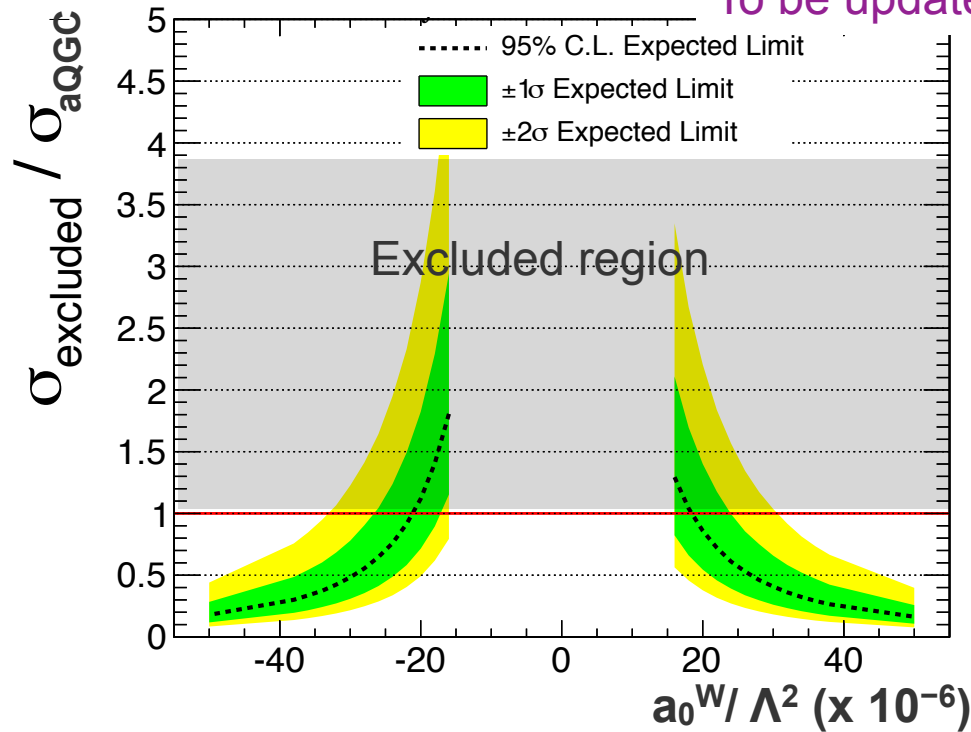
Main systematic uncertainties assumed:

- aQGC signal strength 30%
- Bkgd. normalization 20%
- Experimental uncertainties (JES/R, efficiencies, luminosity,...) each within 5%.

Studies ongoing to understand whether  $m(WW\gamma)$  is equally performant and/or less susceptible to non-unitarity.

# Expected limits on aQGCs

To be updated (used same k-factors for SM & aQGC)



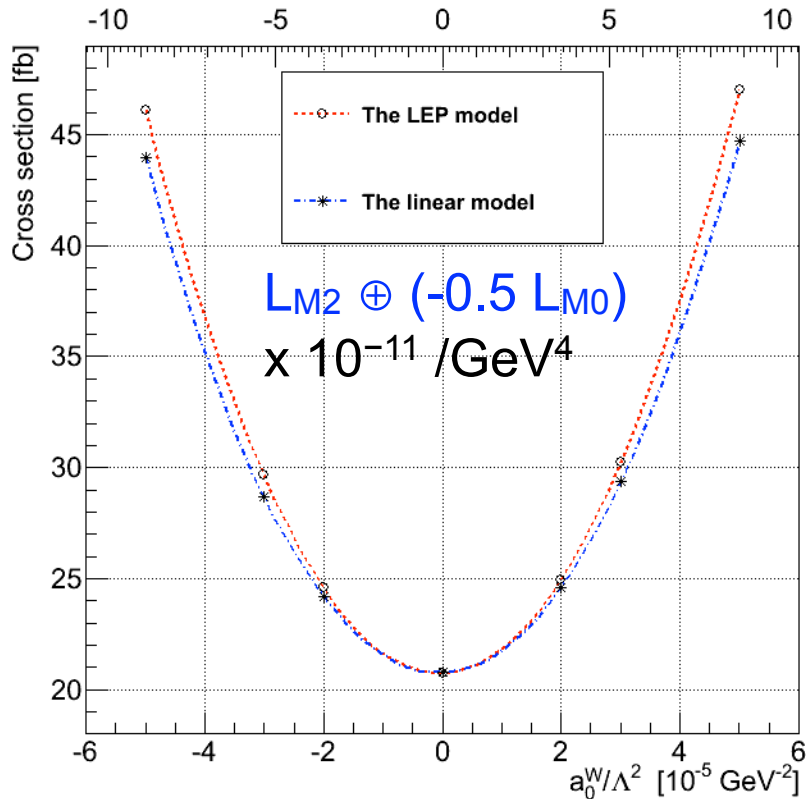
Expected limits:  $-2 \times 10^{-5} < a_0^W / \Lambda^2 < 2 \times 10^{-5}$        $-3 \times 10^{-5} < a_c^W / \Lambda^2 < 3 \times 10^{-5}$

LEP limits  
[NB: uses FF]

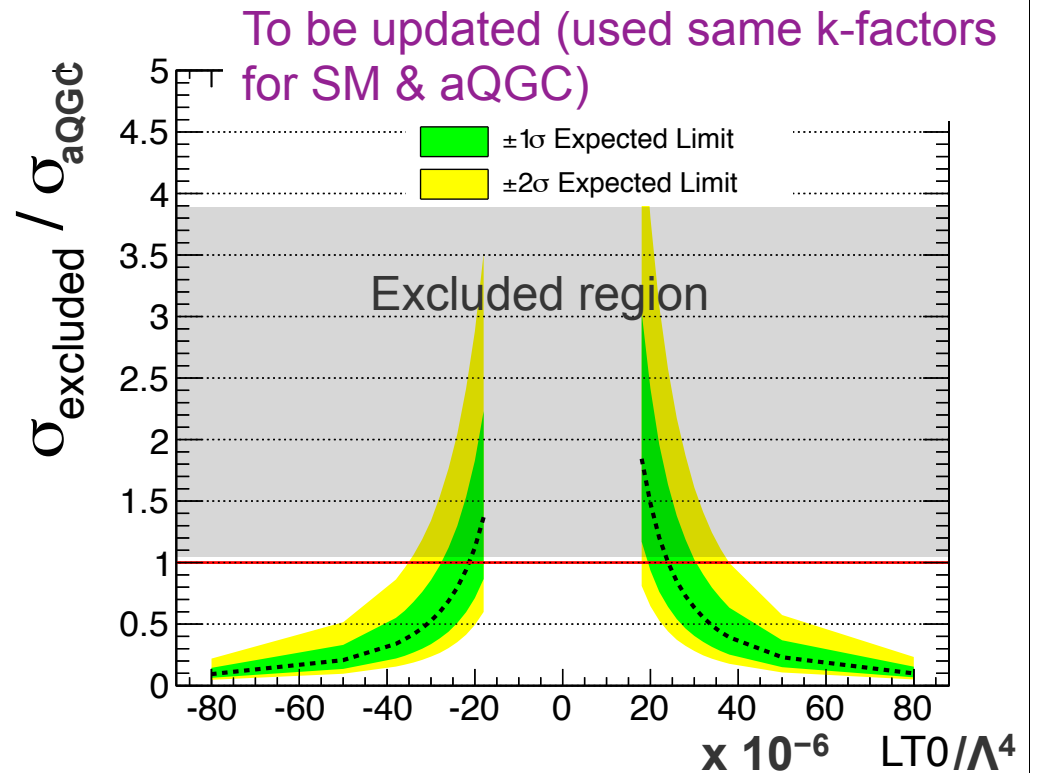
|                |         |  |
|----------------|---------|--|
| $W^+W^-\gamma$ | $a_0^W$ | $-0.020 \text{ GeV}^{-2} < a_0^W / \Lambda^2 < 0.020 \text{ GeV}^{-2}$ |
| $W^+W^-\gamma$ | $a_c^W$ | $-0.053 \text{ GeV}^{-2} < a_c^W / \Lambda^2 < 0.037 \text{ GeV}^{-2}$ |

0(10<sup>2</sup>) times more precise than LEP combined limit.  
Less precise than exclusive  $\gamma\gamma \rightarrow WW$  limit.

# Expected limits II



Shows that coupling  $a_0$  in D6 realization can be expressed as linear combination of couplings  $L_M$  of D8 realization



Couplings  $L_T$  are novel to D8 realization. There is no D6 equivalent. We set limit on  $L_{T0}$  assuming other  $L_T$ 's vanish (they all produce the ~same effect).

# Details of limit setting machinery

- Follow the **RooStat-based framework** used for LHC Higgs combination  
<https://twiki.cern.ch/twiki/bin/view/LHCPhysics/HiggsCombination>
- Description of ATLAS and CMS interfaces to this framework – and their validation of each others results – can be found here  
<http://indico.cern.ch/conferenceDisplay.py?confId=120429>  
(in particular, talks by Kyle Cranmer and Giovanni Petrucciani)

Limits shown in the previous slides computed using CMS interface to RooStats

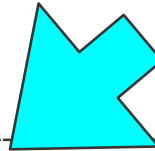
**(Used CLs limit and LHC test statistic, i.e., full profiling of nuisance parameters. See above links for detail.)**

## A binary package wrapping RooStats

- Takes as input either a simple text datacard for counting experiments (same format as L&S) or any RooStats HighLevelFactory file.
- Configures and runs RooStats methods, prints results and saves them to root files.
- Takes care of generating toys for expected limits, or averaging results of multiple runs.

# Example data card

imax 2 number of bins  
 jmax 1 number of processes minus 1  
 kmax 12 number of nuisance parameters



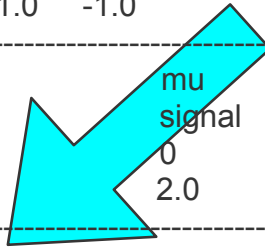
Cleanly tabulated effect on each background due to each source of systematic

|                   |    |             |                                   |
|-------------------|----|-------------|-----------------------------------|
| shapes background | mu | mu_a0W.root | background \$PROCESS_\$SYSTEMATIC |
| shapes data_obs   | mu | mu_a0W.root | data_obs                          |
| shapes signal     | mu | mu_a0W.root | signal_a0w_m2                     |
| shapes background | el | el_a0W.root | background \$PROCESS_\$SYSTEMATIC |
| shapes data_obs   | el | el_a0W.root | data_obs                          |
| shapes signal     | el | el_a0W.root | signal_a0w_m2                     |

|             |      |      |
|-------------|------|------|
| bin         | mu   | el   |
| observation | -1.0 | -1.0 |

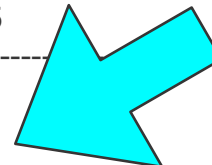
Used lognormal distributions for all systematics

|         |        |            |        |            |
|---------|--------|------------|--------|------------|
| bin     | mu     | el         | mu     | el         |
| process | signal | background | signal | background |
| process | 0      | 1          | 0      | 1          |
| rate    | 2.0    | 35.4       | 1.8    | 38.5       |



Broke systematics down into uncorrelated subsets

|              |        |      |     |      |     |
|--------------|--------|------|-----|------|-----|
| eff_e        | lnN    | -    | -   | 1.03 | -   |
| eff_m        | lnN    | 1.03 | -   | -    | -   |
| JER          | lnN    | 1.02 | -   | 1.02 | -   |
| JES          | lnN    | 1.01 | -   | 1.01 | -   |
| MET          | lnN    | 1.01 | -   | 1.01 | -   |
| PDF          | lnN    | 1.05 | -   | 1.05 | -   |
| signal_norm  | lnN    | 1.30 | -   | 1.30 | -   |
| el_backshape | shape1 | -    | -   | -    | 1.0 |
| lumi         | lnN    | 1.03 | -   | 1.03 | -   |
| mu_backshape | shape1 | -    | 1.0 | -    | -   |
| pileup       | lnN    | 1.01 | -   | 1.01 | -   |



Start with a txt input, define a mathematical representation, and then prepare the ROOT workspace to be fed to the limit setter

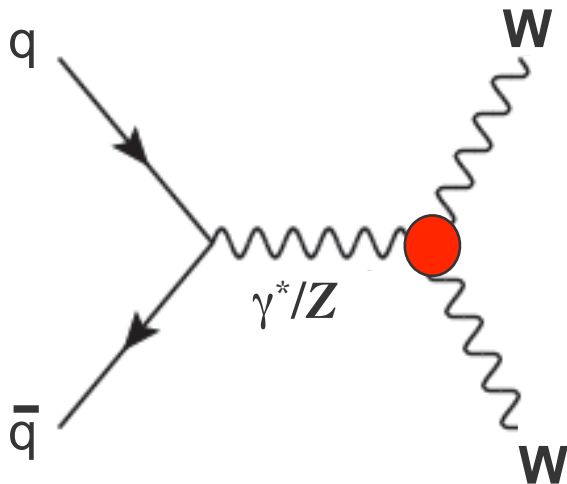
# Summary

- ☑ Study of QGC and related states is a rich physics program
  - LHC data sufficient for sensitivity to SM QGC and aQGCs
  - New excitement after the discovery of a light Higgs boson
- ☑ LHC collaborations have dedicated effort to measure QGCs
  - in both multi-boson and scattering topologies
- ☑ Starting to set serious constraints on EWK gauge boson coupling
  - Already broke new ground with 8 TeV data by exceeding LEP aQGC limits by orders of magnitude
  - More studies underway – including development of new techniques & NLO computation – to improve precision at 13 TeV

**Thank You !**

BACKUP SLIDES

# Measurements of gauge boson self couplings



- Gauge boson trilinear & quartic couplings emerges naturally from the non-abelian gauge symmetry structure of the SM.
- With  $\mathcal{O}(10^3)$  WW,  $\mathcal{O}(10^2)$  WZ, and  $\mathcal{O}(\text{dozens})$  ZZ events, quickly approaching precision measurement of gauge couplings.
  - Already improved over LEP and Tevatron in most cases.
- Measure anomalous coupling parameters in effective Lagrangian approach.

**Let's do a quick overview of the current aTGC results**

in the notation of LEP parametrization [hep-ph/9601233](https://arxiv.org/abs/hep-ph/9601233)

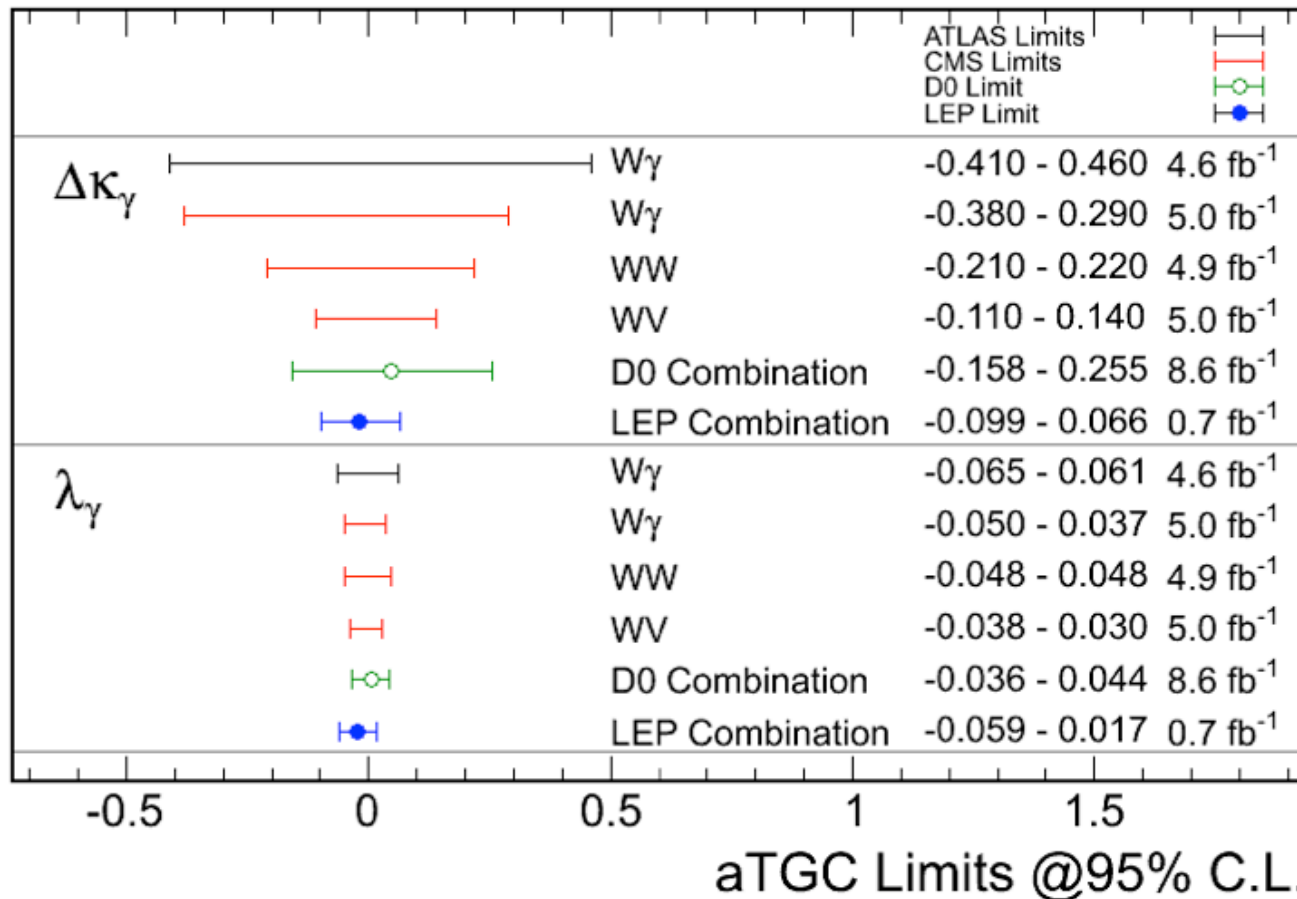
since they are also relevant for discussion of quartic couplings

# Summary of aTGC measurements I

<https://twiki.cern.ch/twiki/bin/view/CMSPublic/PhysicsResultsSMPaTGC>

## Limits on $WW_\gamma$ couplings

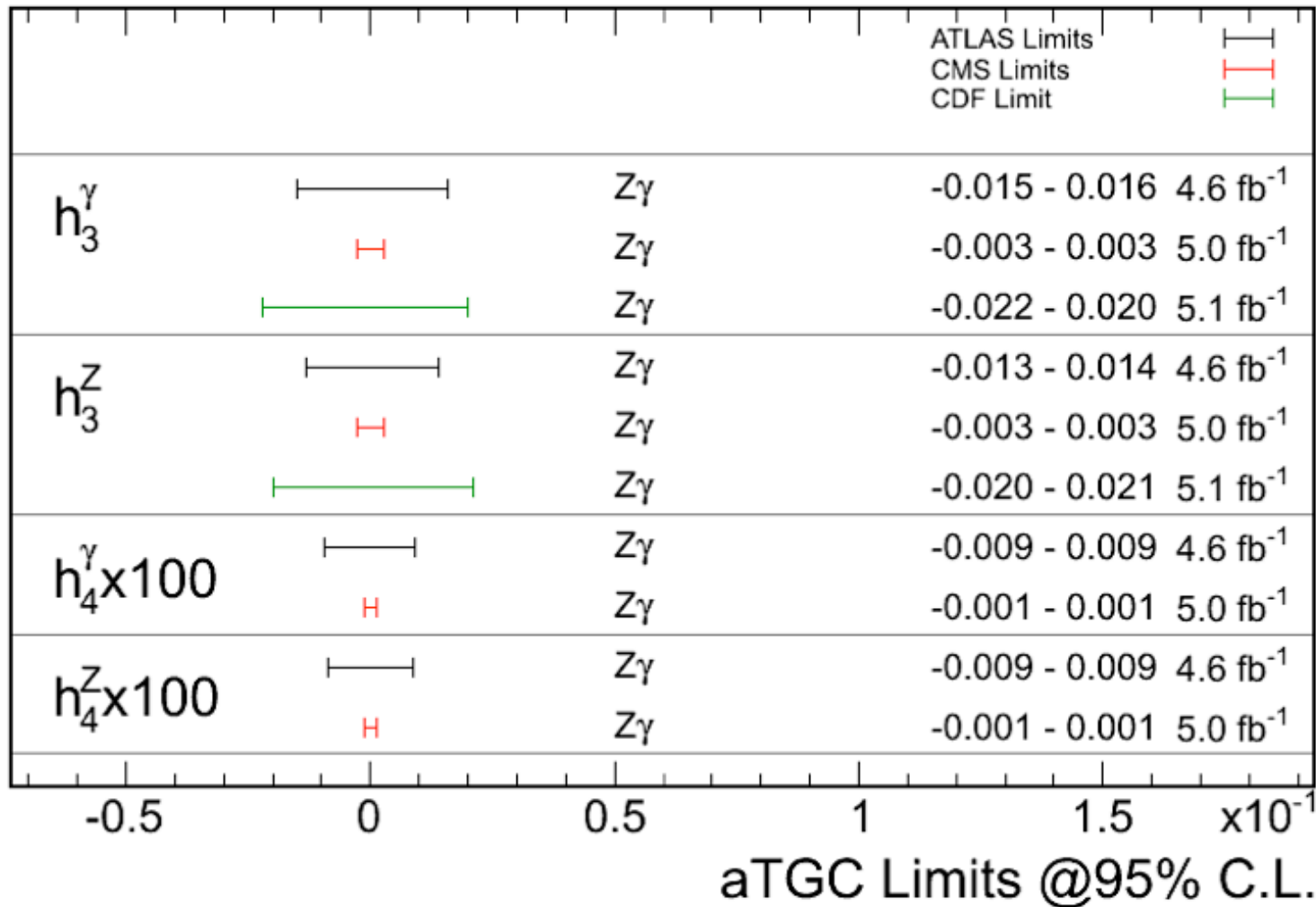
Feb 2013



# Summary of aTGC measurements II

Feb 2013

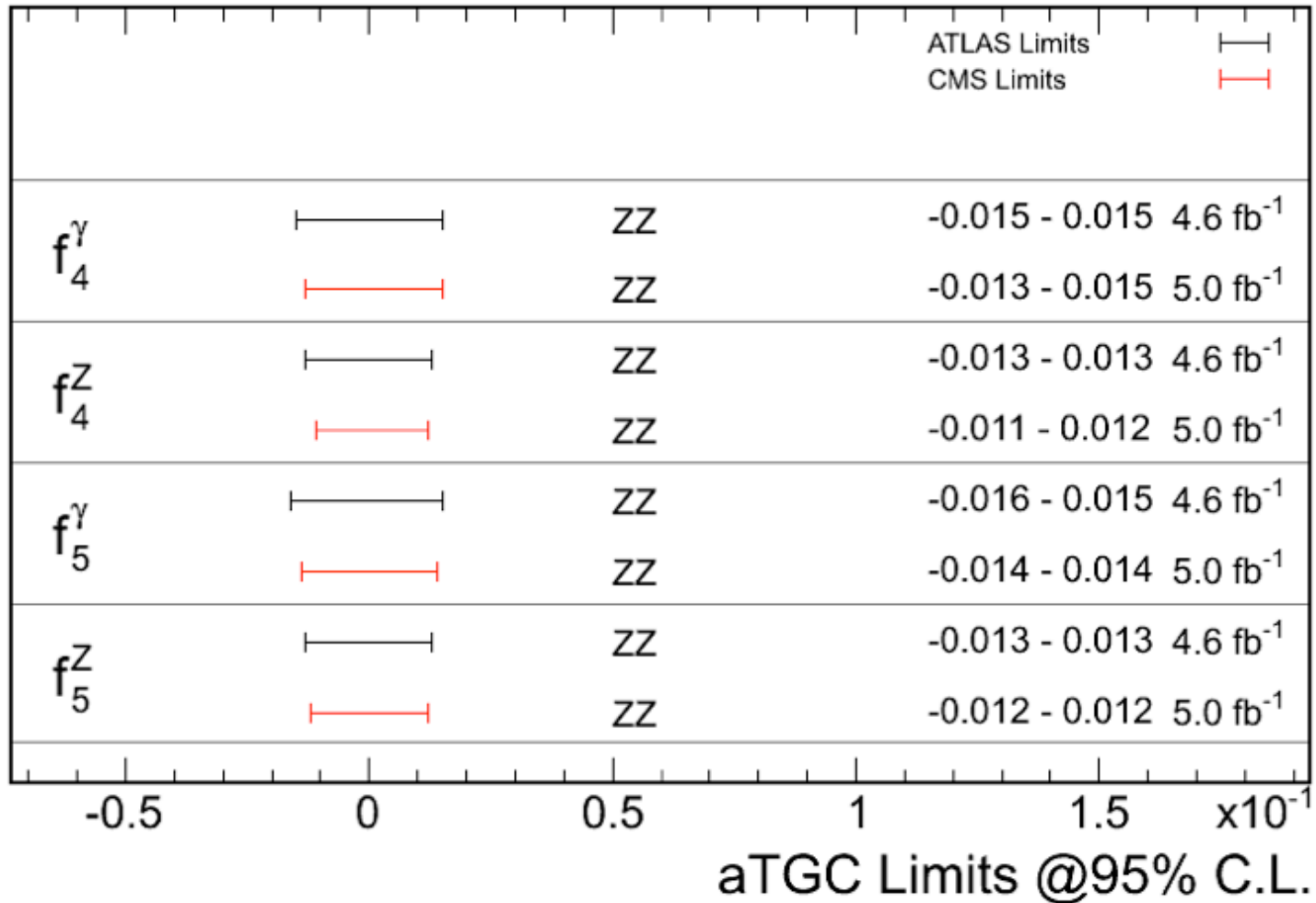
## Limits on $Z\gamma\gamma$ and $ZZ\gamma$ couplings



# Summary of aTGC measurements III

## Limits on $ZZ\gamma$ and $ZZZ$ couplings

Feb 2013

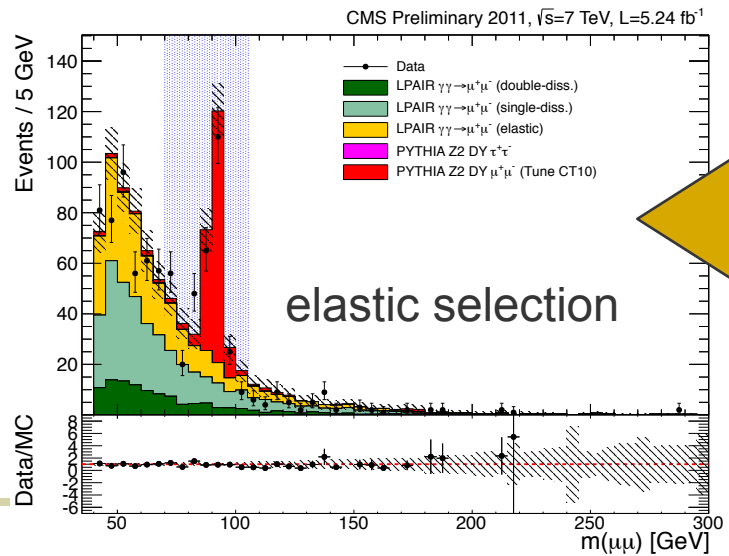


# $\gamma\gamma \rightarrow WW$ : CMS analysis details

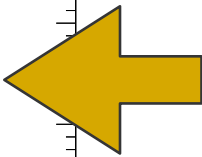
## Event Selection:

- lepton  $p_T > 20$  GeV,  $|\eta| < 2.4$ , isolated and well-identified
- $m(\mu^\pm e^\mp) > 20$  GeV,  $p_T(\mu^\pm e^\mp) > 30$  GeV (to reduce  $\gamma\gamma \rightarrow \tau^+\tau^-$ )
- No extra tracks associated with  $\mu^\pm e^\mp$  vertex

| Selection step                             | Signal $\epsilon \times A$ | Visible cross section (fb) | Events in data |
|--|----------------------------|----------------------------|----------------|
| Trigger and preselection                   | 28.5%                      | 1.4                        | 9086           |
| $m(\mu^\pm e^\mp) > 20$ GeV                | 28.0%                      | 1.4                        | 8200           |
| Muon ID and Electron ID                    | 22.6%                      | 1.1                        | 1222           |
| $\mu^\pm e^\mp$ vertex with 0 extra tracks | 13.7%                      | 0.7                        | 6              |
| $p_T(\mu^\pm e^\mp) > 30$ GeV              | 10.6%                      | 0.5                        | 2              |

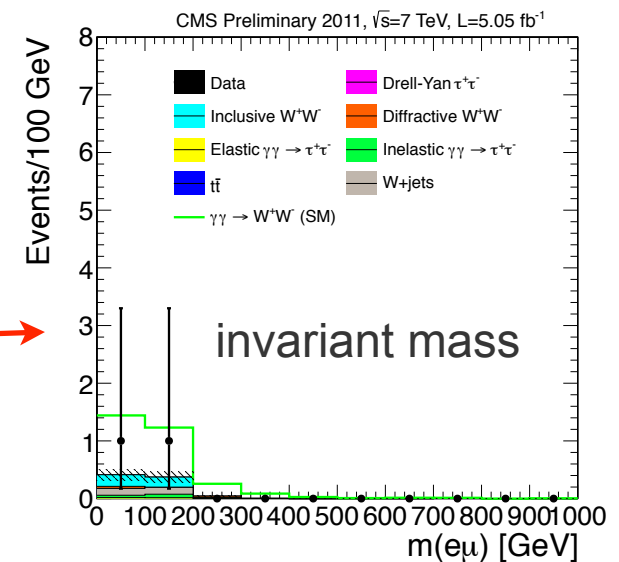
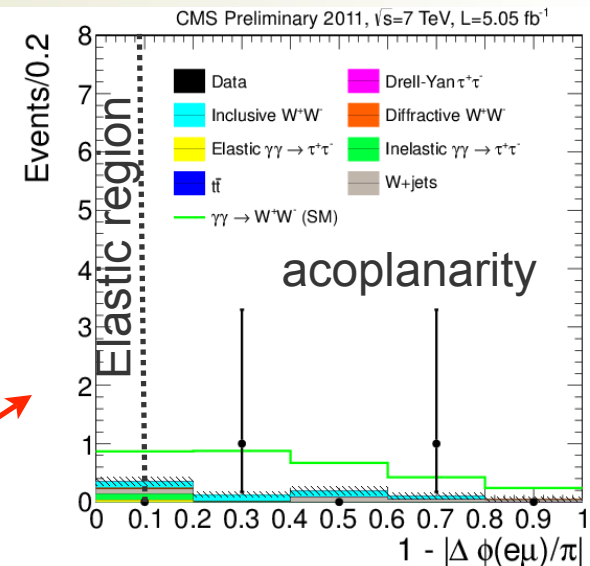
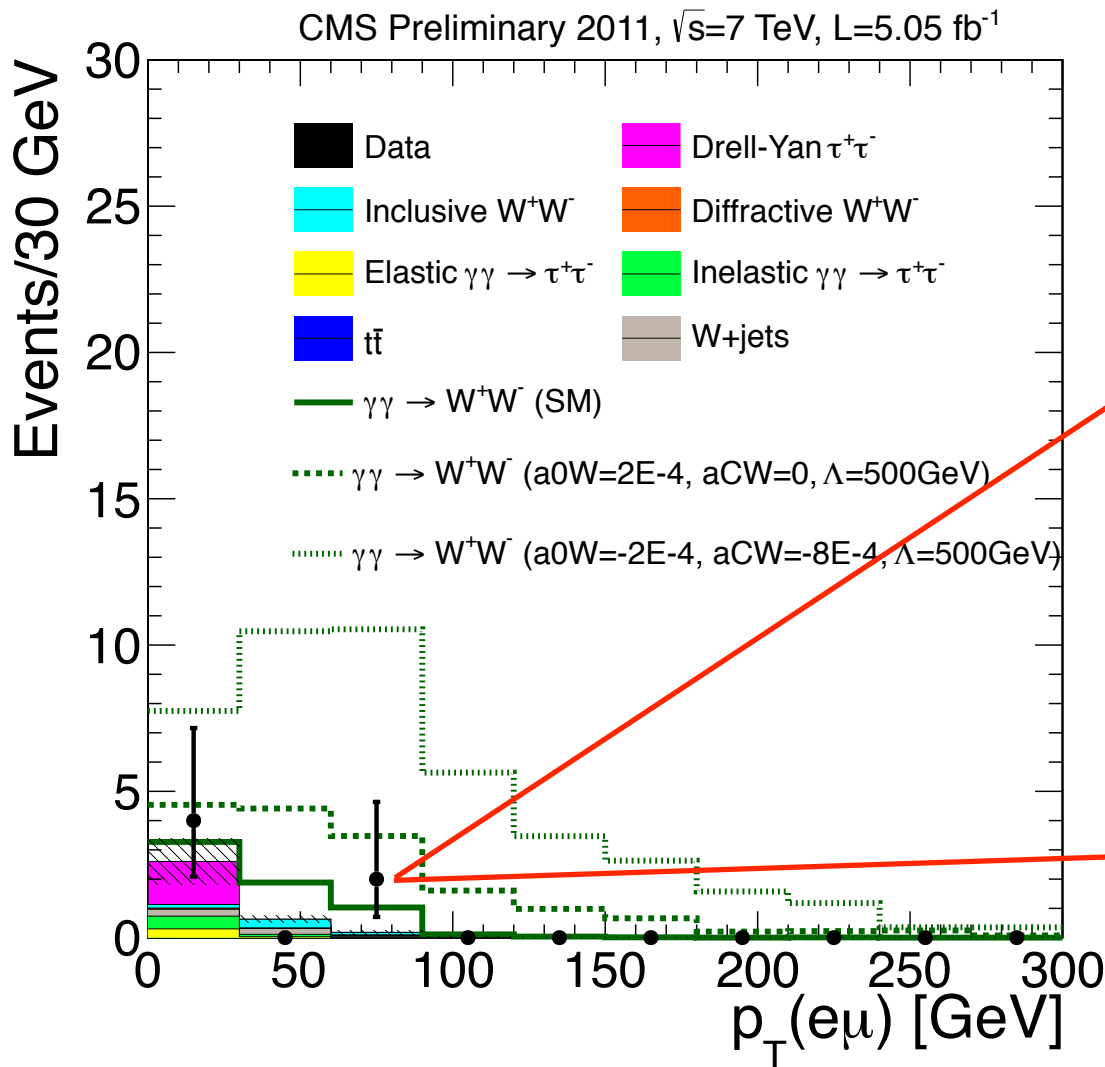


(Expect  $2.2 \pm 0.5$  signal,  $0.84 \pm 0.13$  bkgd)



Use exclusive  $\mu^+\mu^-$  production as benchmark to validate efficiency of vertexing and exclusivity selection and pileup dependence.

# Kinematic distributions of signal-like events



# Limits on aQGC

Observe no events in the high  $p_T$  region where SM contribution is small within the acceptance of  $p_T(\mu, e) > 20 \text{ GeV}$ ,  $|\eta(\mu, e)| < 2.4$ ,  $p_T(\mu^\pm e^\mp) > 100 \text{ GeV}$ :  

$$\sigma(pp \rightarrow p^{(*)}W^+W^-p^{(*)} \rightarrow p^{(*)}\mu^\pm e^\mp p^{(*)}) < 1.9 \text{ fb.}$$

## Limits on aQGC without form-factors (LHC preferred way):

$$-2.80 \times 10^{-6} < a_0^W / \Lambda^2 < 2.80 \times 10^{-6} \text{ GeV}^{-2} \quad (a_C^W / \Lambda^2 = 0, \text{ no form factor}),$$

$$-1.02 \times 10^{-5} < a_C^W / \Lambda^2 < 1.02 \times 10^{-5} \text{ GeV}^{-2} \quad (a_0^W / \Lambda^2 = 0, \text{ no form factor}),$$

## Limits using a form-factor:

$$-0.00017 < a_0^W / \Lambda^2 < 0.00017 \text{ GeV}^{-2} \quad (a_C^W / \Lambda^2 = 0, \Lambda = 500 \text{ GeV}),$$

$$-0.0006 < a_C^W / \Lambda^2 < 0.0006 \text{ GeV}^{-2} \quad (a_0^W / \Lambda^2 = 0, \Lambda = 500 \text{ GeV}),$$

where the dipole form factor is

$$a_{0,C}^W(W_{\gamma\gamma}^2) = \frac{a_{0,C}^W}{\left(1 + \frac{W_{\gamma\gamma}^2}{\Lambda^2}\right)^p}$$

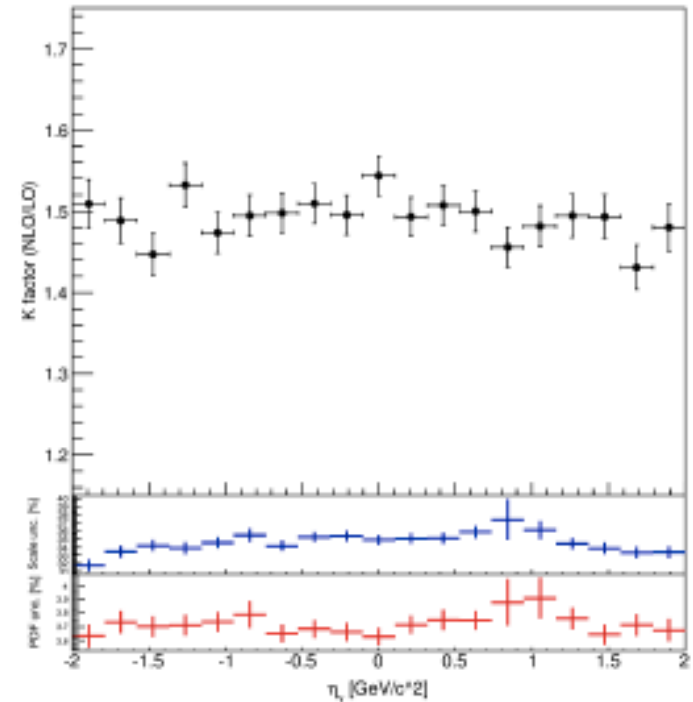
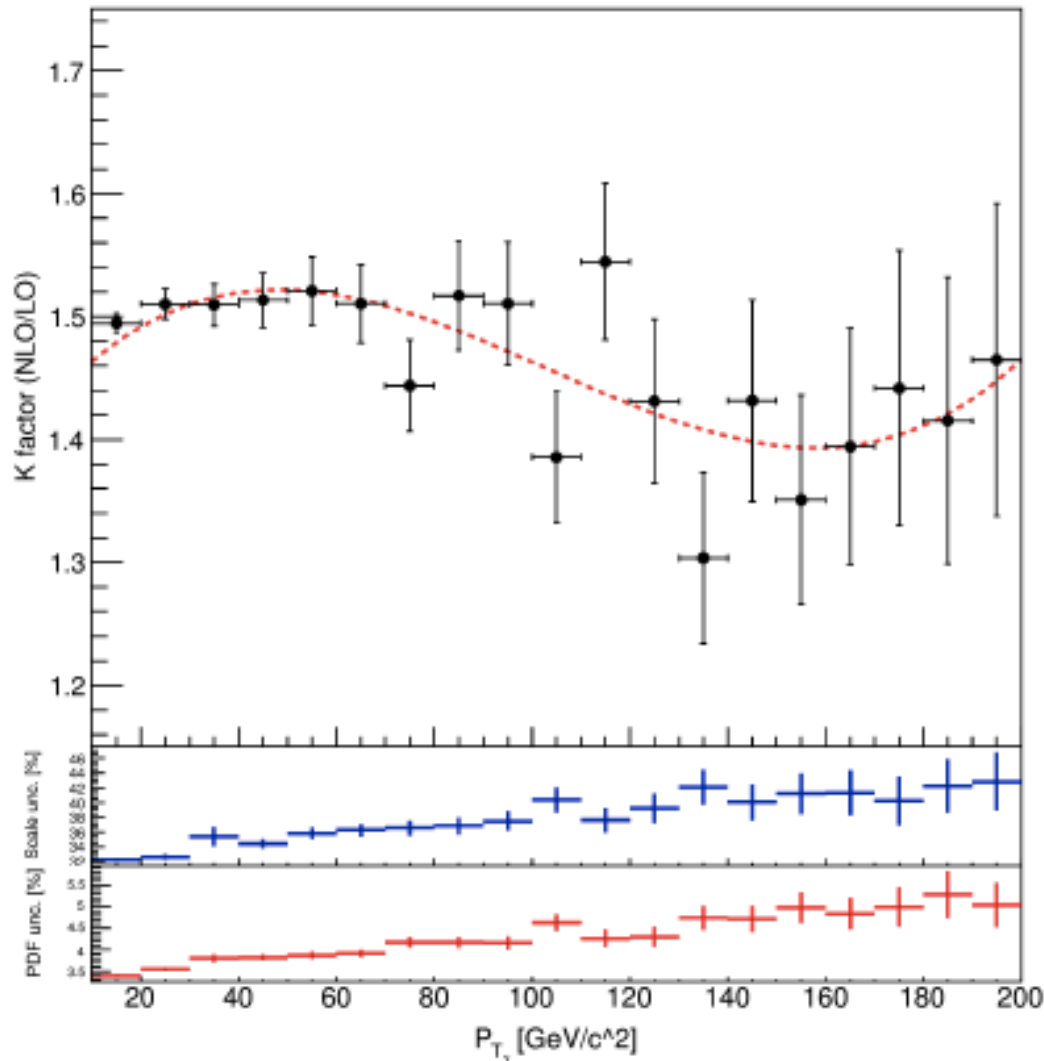
$W_{\gamma\gamma} = \gamma\gamma$  center of mass energy

$p =$  a free parameter = 2 by convention

Two orders of magnitude more constraining than the LEP combined limit.

# SM $WW\gamma$ k-factor after requiring jet veto

Additional jet veto for  $p_T > 30\text{GeV}$  and  $|\eta| < 4.5$



Clearly, applying jet veto in this analysis is not a good idea !!!

## Computation of scale and PDF uncertainties

- Reweight to get scale dependence and PDF uncertainty Ref: arXiv1110.4738

### Scale uncertainty

Factor 0.5/2 around central scale

3x3 values of weight

$\delta S = \max - \min$

### PDF uncertainty

MSTW2008nlo68cl

1 central + 20 pairs

arXiv: 0201195v3 Eq(3)

$$\Delta X = \frac{1}{2} \left( \sum_{i=1}^{N_p} [X(S_i^+) - X(S_i^-)]^2 \right)^{1/2}$$